



## Magnetic properties of $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$ ( $x=0.0-0.5$ ) bulk metallic glasses

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### ABSTRACT

The effects of partial substitution of Fe by Ni in  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4$ , and  $0.5$ ) bulk metallic glasses (BMGs) on their magnetic properties were studied. It is found that the saturation polarization decreases from 1.15 T to 0.69 T with increasing Ni content from  $x=0.0$  to  $x=0.5$ ; the Curie temperature reaches its maximum of 598 K for composition  $x=0.1$ , and is then followed by a quick decrease with higher Ni content. Moreover, the random anisotropy and mean field theory were developed to investigate the magnetic properties of these BMGs, and the variations of saturation polarization and Curie temperature were well explained by calculating the magnetic exchange stiffness constant  $A$  and the nearest neighbor transition-metal-pair exchange interactions  $J$  by using the model.

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### 1. Introduction

Bulk metallic glasses (BMGs) have attracted great attention due to their unique properties and potential applications [1–3]. Among them, the Fe-based BMGs are more attractive for application since they not only possess many good properties, such as excellent soft magnetic properties [4,5], besides ultrahigh strength [6–8] and good anticorrosion property [9], but are also much cheaper in comparison to other BMGs [10]. Recently, Fe-based BMGs have been intensively investigated and their critical size has been increased [11–16]. It is found that partial substitution of Fe with other magnetic elements Ni or Co may remarkably enhance the GFA [17,18] and change the soft magnetic properties due to compositional variation [19,20]; nevertheless, it is also found that the effects of Ni and Co on magnetic properties are somewhat different from each other [21,22], even if they are at neighboring positions in the Element Periodic Table. However, little attention concerning how and why the magnetic properties change induced by partial substitution of Fe by Ni was paid, and even less detailed theoretical research on the mechanism of partial substitution of Fe with constituting elements in Fe-based BMGs was carried out.

In the present work, we investigated the effects of partial substitution of Fe by Ni on the thermal characteristics and

magnetic properties of Fe–B–Si–Nb BMGs considering its high GFA and excellent magnetic properties [12,18]. On this basis, the random anisotropy and mean field theory of the magnetic properties for  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4$ , and  $0.5$ ) were developed to explain the variations of saturation polarization and Curie temperature.

### 2. Experiments

Multi-component  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4$ , and  $0.5$ ) ingots were prepared by arc melting the mixtures of Fe (99.99%), Ni (99.99%) and Nb (99.99%) metals, and B (99.5%) and Si (99.999%) crystals in an argon atmosphere. Cylindrical alloy rods with diameters of 1.5 mm were produced by the copper mold casting method. The structures of as-cast samples were identified by X-ray diffraction (XRD) with  $\text{Cu K}\alpha$  radiation. The thermal stability of the glassy samples was examined using a NETZSCH 404 C differential scanning calorimeter (DSC) at a heating rate of 0.67 K/s under a flow of high purity argon. The Curie temperature of the as-cast alloys was measured by a physical property measurement system (PPMS) at a heating rate of 0.167 K/s. For magnetic properties investigation,  $M$ – $H$  hysteresis loops were measured with a vibrating sample magnetometer (VSM) at ambient temperatures. Additionally, the density of the specimens was measured using Archimedes's method with the accuracy less than 1%.

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### 3. Results and discussion

Fig. 1 shows the XRD patterns for all six as-cast rods with diameter of 1.5 mm. It can be seen that each pattern contains only broad maxima without any additional sharp Bragg peaks, and the diffraction images did not show any traces of crystallinity. This is characteristic of fully amorphous samples. It is found that as the content of Ni increases, the diffraction angle of the principal diffuse peak in XRD pattern is almost unchanged. The position of an X-ray halo maximum is known to be directly related to the average radius of the first coordination shell,  $r_1$ , according to the Bragg equation:  $2r_1 \sin \theta = \lambda$  with  $\lambda$  the X-ray wavelength and  $2\theta$  the scatter angle corresponding to the halo maximum. The almost unchanged first peak might be due to the almost equal radius of Ni (0.125 nm) and Fe (0.124 nm) [23]; even if Fe was replaced by Ni, the value of  $r_1$  remains almost unchanged, which leads to nearly invariant  $2\theta$ .

Fig. 2(a) presents the DSC traces of the as-cast  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  BMGs ( $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) measured at a constant heating rate of 0.67 K/s. It can be seen that all DSCs exhibit one endothermic event characteristic of the glass transition from the amorphous solid to the supercooled, followed by a crystallization exothermic reaction, which further confirms the glassy nature of the rods. Although  $T_g$  and  $T_x$  decrease from 837 to 776 K and from 882 to 834 K, respectively,  $\Delta T_x$  increases gradually from 45 to 58 K with increasing Ni content to  $x=0.5$ . Thus, it is suggested that the thermal stability of the supercooled liquid increases with increasing Ni. Fig. 2(b) describes DSC curves revealing the cooling behavior of the same Fe-based BMGs system, which demonstrates that the  $\text{Fe}_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  BMG exhibits the highest  $T_l$  and  $T_e$  of 1479 and 1405 K, respectively, as well as the largest temperature interval ( $T_l - T_e$ ) of 74 K, implying that this composition does not lie in the vicinity of a eutectic point. With increasing Ni content from  $x=0.1$  to  $x=0.5$ , the  $T_l$  and  $T_e$  decrease gradually from 1471 to 1298 K and from 1405 to 1291 K, respectively, indicating that the alloy approaches a eutectic point in this process; especially when  $x=0.4$ ,  $T_l$  is very close to  $T_e$ ; thus, it is considered to be the closest to the eutectic point in this alloy system.

#### 3.1. Saturation polarization

As expected, the fully glassy samples display excellent magnetic properties. Table 1 summarizes the saturation magnetization  $M_s$  (given in  $\text{Am}^2/\text{kg}$  and in Bohr magnetons  $\mu_B$  per magnetic atom), saturation polarization  $J_s$  and Curie temperature  $T_c$  as well as the density  $\rho$  for as-cast samples. Here, the saturation

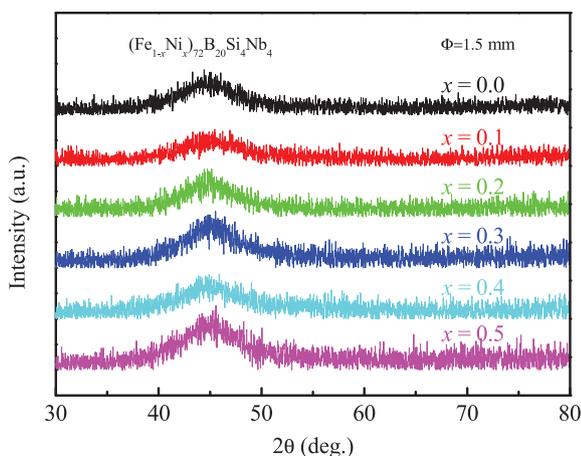


Fig. 1. XRD patterns of the cast  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) alloys rods with diameters of 1.5 mm.

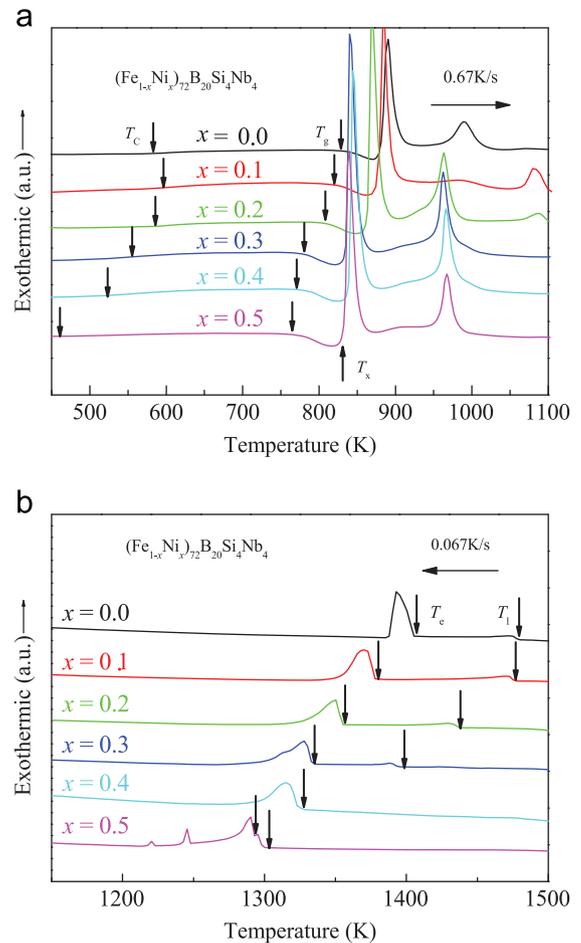


Fig. 2. (a) DSC traces measured at a constant heating rate of 0.67 K/s for cast  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) alloys. (b) DSC traces measured at a constant cooling rate of 0.067 K/s for cast  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) alloys.

Table 1

Magnetic properties data for  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  glassy rods: magnetization at saturation  $M_s$  (given in  $\text{A m}^2/\text{kg}$  and in Bohr magnetons  $\mu_B$  per magnetic atom), saturation polarization  $J_s$ , and Curie temperature  $T_c$ . Additionally, the density  $\rho$  is also shown.

Atomic fraction $x$	$M_s$ ( $\text{A m}^2/\text{kg}$ )	$J_s$ (T)	$M_s$ ( $\mu_B$ )	$T_c$ (K)	$\rho$ ( $\text{g}/\text{cm}^3$ )
0.0	123.2	1.15	1.25	581	7.4377
0.1	115.4	1.08	1.14	598	7.4493
0.2	104.9	0.99	1.05	588	7.4868
0.3	95.5	0.91	0.95	565	7.5794
0.4	81.6	0.78	0.81	525	7.6559
0.5	70.8	0.69	0.71	458	7.7576

polarization was calculated by using the actual density values ( $J_s = 4\pi \times 10^{-7} \rho M_s$ ), as presented in Fig. 3, and the saturation polarization as a function of Ni content is shown in Fig. 4. We can see that the saturation polarization decreases monotonically from 1.15 to 0.69 T upon increasing Ni content from  $x=0.0$  to 0.5.

The variation of saturation polarization as a function of Ni content should be explained in more detail. From the random anisotropy theory [24], the correlation between the saturation polarization and the exchange stiffness constant  $A$  can be expressed as

$$J_s = \frac{8 \times 10^{-7} \pi \rho A}{H_{ex} R_a^2} \quad (1)$$

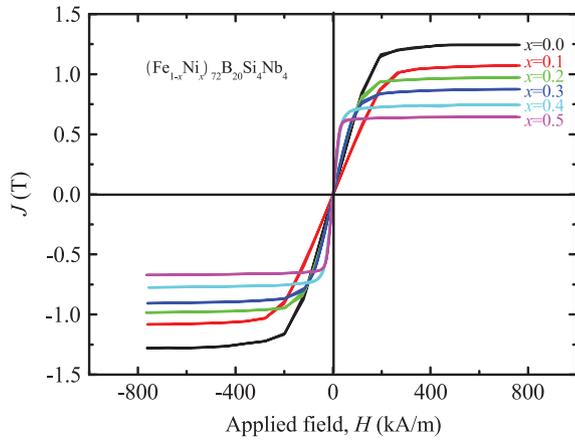


Fig. 3. DC hysteresis loops for cast  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) glassy samples.

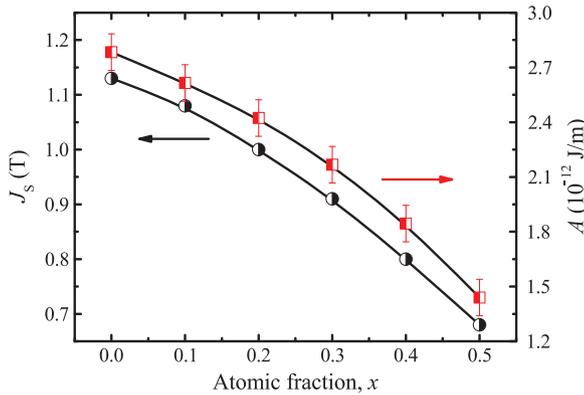


Fig. 4. Variations of saturation polarization and exchange stiffness constant  $A$  as a function of the atomic fraction  $x$ .

where  $H_{ex}$  is the exchange field which remains unchanged,  $R_d$  the length over which the local axes show a correlation (typically  $10 \text{ \AA}$ ) [25], and  $\rho$  the density. Consequently, the saturation polarization  $J_s$  is proportional to the exchange stiffness constant  $A$ , and the latter can be obtained from the Curie temperature  $T_C$ . On the basis of Hasegawa's model, one can derive the following equation [26]:

$$A = \frac{x_{\text{Fe}} S_{\text{Fe-Fe}} k_B T_C}{4(S_{\text{Fe-Fe}} + 1)r_{\text{Fe-Fe}}} \quad (2)$$

where  $x_{\text{Fe}}$  is the concentration of Fe in atomic percent,  $S_{\text{Fe-Fe}}$  the spin moment of Fe,  $T_C$  the Curie temperature,  $k_B$  the Boltzmann constant, and  $r_{\text{Fe-Fe}}$  the nearest-neighbor distance, which is considered to be twice the Fe atomic radius. Moreover,  $S_{\text{Fe-Fe}}$  can be calculated by

$$S_{\text{Fe-Fe}} = M(\mu_B) / 2x_{\text{Fe}} \mu_B \quad (3)$$

in which  $M(\mu_B)$  is the magnetic moment of the samples considering Bohr magnetons ( $M(\mu_B) = JSZ / \rho N_A$ , and here  $Z$  is the molar mass). However, Eq. (3) is valid only when the starting alloy sample ( $x=0.0$ ) or the Fe-Ni and Ni-Ni interactions could be ignored.

Normally, when Fe is partially substituted by Ni, the magnetic saturation should be considered as

$$M(\mu_B) = wM_{\text{Fe}}(\mu_B) + (1-w)M_{\text{Ni}}(\mu_B) \quad (4)$$

in which  $M_{\text{Fe}}(\mu_B)$  and  $M_{\text{Ni}}(\mu_B)$  are the magnetic moments of Ni-free and Fe-free amorphous alloys respectively, and  $w$  is the atomic fraction  $\text{Fe}/(\text{Fe}+\text{Ni})$ .

From the above analysis, the exchange stiffness constant  $A$  as a function of the atomic fraction  $x$  (the Ni content) was calculated,

as shown in Fig. 4. It can be seen that the exchange stiffness constant  $A$  decreases from  $2.73 \times 10^{-12} \text{ J/m}$  to  $1.47 \times 10^{-12} \text{ J/m}$  when  $x$  increases from 0 to 0.5. The decrease of the stiffness constant means the promotion of the magnetic anisotropy and accordingly a decrease of exchange energy; as a result, the saturation magnetization decreases accordingly with increasing Ni content.

### 3.2. Curie temperature

Values of the  $T_C$  for various compositions of  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  BMGs are shown in Fig. 5. Fig. 5 presents a number of salient  $T_C$  going through a broad maximum at  $x=0.1$  for the alloy series, which decreases rapidly as  $x$  increases. The corresponding values are also summarized in Table 1. Moreover, in order to confirm this phenomenon, the thermomagnetic curves, measured at a constant heating rate of  $0.167 \text{ K/s}$  up to  $800 \text{ K}$ , are presented in Fig. 6(a). Apparently, all the samples did not start to crystallize yet and the values correspond to the glassy states. The saturation magnetization can be described by  $M_S(T) = M(0)(1 - T/T_C)^{0.36}$  [20,27]. In order to minimize the errors, the experimental results were plotted as  $M_S(T)^{0.36}$  vs  $T$ , shown in Fig. 6(b). It was considered that  $T_C$  deviates from linearity. We can see that the  $T_C$  measured by DSC is in agreement with that obtained from PPMS. The small deviation can be attributed to the experimental error. The  $T_C$  goes through a broad maximum at  $x=0.1$  for Fe-rich alloys, indicating a large Fe-Fe exchange ( $J_{\text{Fe-Fe}}$ ) which is ferromagnetic, whereas the sharp decrease in the values of  $T_C$  for the Ni-rich alloys suggests that the Ni-Ni exchange ( $J_{\text{Ni-Ni}}$ ) tends to be weaker and weaker as Ni concentration increases. In order to arrive at a quantitative estimation of these exchange interactions, it is imperative to compare the experimentally observed dependence of  $T_C$  on  $x$  with theoretical models. It is obvious that the mean-field model can be used to explain the variation of  $T_C$  with  $x$  [28]. According to this theory, the calculation of  $T_C$  for two sublattices reduces to

$$T_C = \frac{1}{2} [T_{\text{Ni-Ni}}x + T_{\text{Fe-Fe}}(1-x)] + \left\{ \frac{1}{4} [T_{\text{Ni-Ni}}x + T_{\text{Fe-Fe}}(1-x^2)]^2 + T_{\text{Fe-Ni}}^2 x(1-x) \right\}^{1/2} \quad (5)$$

where

$$T_{\text{Fe-Fe}} = \frac{S_{\text{Fe-Fe}}(S_{\text{Fe-Fe}} + 1)J_{\text{Fe-Fe}}}{3k_B}$$

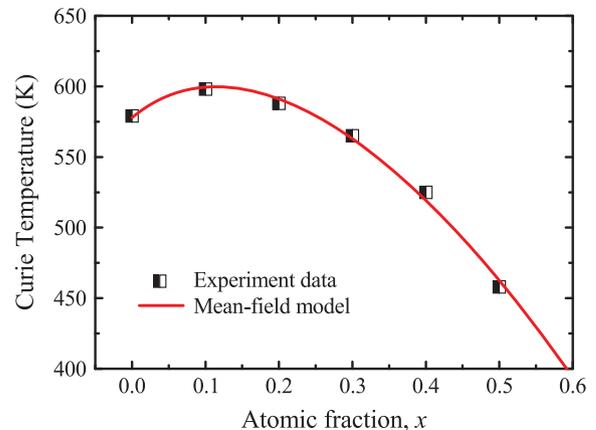
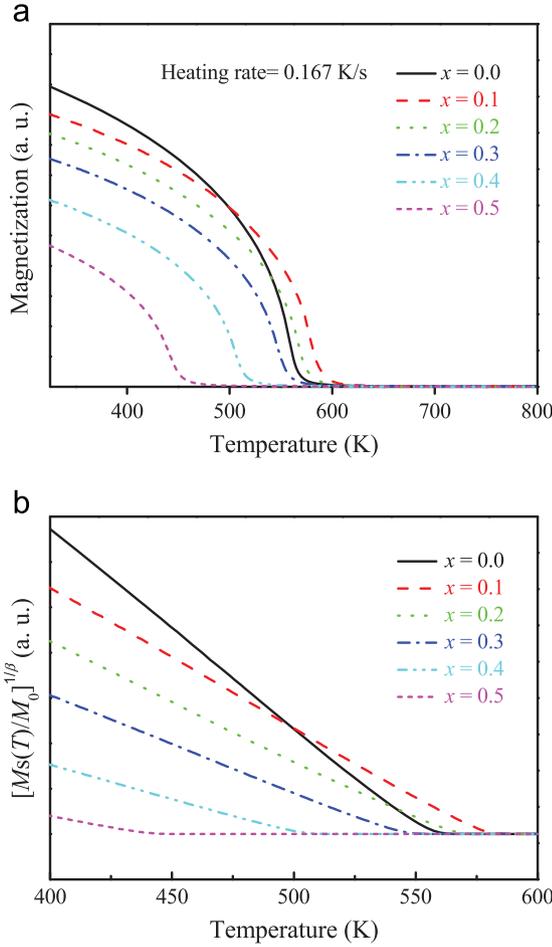


Fig. 5. Comparisons of experimental and theoretical data of Curie temperature for  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  as-cast glassy samples with the atomic fraction  $x$ .



**Fig. 6.** (a) Thermomagnetic curves for as-cast  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) glassy samples. (b) The relationship between  $M_s(T)^{0.36}$  and  $T$  for as-cast  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) glassy samples.

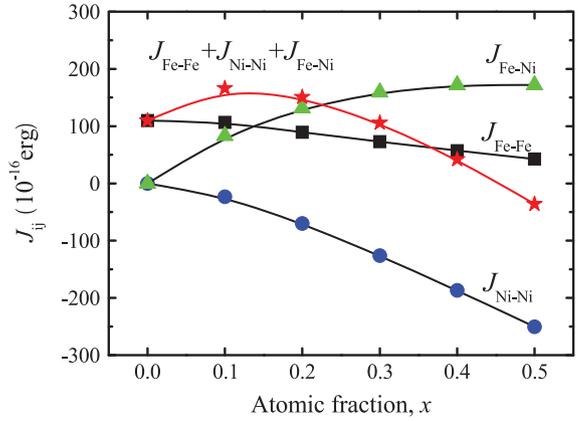
$$T_{\text{Fe-Ni}} = \frac{\sqrt{S_{\text{Ni-Ni}}S_{\text{Fe-Fe}}(S_{\text{Ni-Ni}}+1)(S_{\text{Fe-Fe}}+1)}J_{\text{Fe-Ni}}}{3k_B}$$

$$T_{\text{Ni-Ni}} = \frac{S_{\text{Ni-Ni}}(S_{\text{Ni-Ni}}+1)J_{\text{Ni-Ni}}}{3k_B}$$

where  $J_{ij}$  ( $i=\text{Fe}, j=\text{Ni}$ ) is the exchange interaction. The experimental and corresponding theoretical calculations of  $T_C$  for  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  ( $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$ ) alloys are shown in Fig. 5. We can clearly see that the fitting curves from our model are in good agreement with the experiments. Here, the parameters  $T_{\text{Fe-Fe}}, T_{\text{Ni-Ni}},$  and  $T_{\text{Fe-Ni}}$  obtained from the fitting of experimental data by using Eq. (5) are 575 K, -318 K, and 754 K, respectively, which are in agreement with Ref. [29]. Moreover, the variation of the exchange interactions  $J_{ij}$  as a function of the atomic fraction  $x$  was also obtained, as shown in Fig. 7. It reveals that magnitude of both  $J_{\text{Fe-Fe}}$  and  $J_{\text{Ni-Ni}}$  decreases with increasing Ni concentration; however, the magnitude of  $J_{\text{Fe-Ni}}$  increases with increasing Ni. The total nearest neighbor transition-metal-pair exchange interactions can be expressed as

$$J = J_{\text{Fe-Fe}} + J_{\text{Fe-Ni}} + J_{\text{Ni-Ni}} \quad (6)$$

Thus, we can get that  $J$  goes through a broad maximum at  $x=0.1$ , and then decreases sharply as  $x$  increases, just as shown in Fig. 7, which demonstrates that the results can be explained by the present work.



**Fig. 7.** Variation of the exchange interactions  $J_{ij}$  as a function of the atomic fraction  $x$ .

#### 4. Conclusions

In this paper, the effects of substitution of Fe by Ni on magnetic properties of  $(\text{Fe}_{1-x}\text{Ni}_x)_{72}\text{B}_{20}\text{Si}_4\text{Nb}_4$  BMGs with  $x=0.0, 0.1, 0.2, 0.3, 0.4,$  and  $0.5$  were studied. It is drawn that partial substitution of Fe by Ni results in the decreased saturation polarization and increased Curie temperature, but first goes through a maximum for  $x=0.1$ . The exchange stiffness constant calculated by the random anisotropy model follows the trend of experiments, and the exchange interaction shows a maximum at  $x=0.1$  for Ni in the alloy from mean field theory.

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